

Simple formulas for constellations and bipartite maps with prescribed degrees

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Abstract. We obtain simple quadratic recurrence formulas counting bipartite maps on surfaces with prescribed degrees (in particular, $2k$ -angulations), and constellations. These formulas are the fastest known way of computing these numbers.

Our work is a natural extension of previous works on integrable hierarchies (2-Toda and KP), namely the Pandharipande recursion for Hurwitz numbers (proven by Okounkov and simplified by Dubrovin–Yang–Zagier), as well as formulas for several models of maps (Goulden–Jackson, Carrell–Chapuy, Kazarian–Zograf). As for those formulas, a bijective interpretation is still to be found.

We also include a formula for monotone simple Hurwitz numbers derived in the same fashion.

Keywords: maps, Hurwitz numbers, Toda hierarchy, constellations

1 Introduction

A map is a combinatorial object describing the embedding up to homeomorphism of a multigraph on a compact oriented surface. A bipartite map is a map with black and white vertices, each edge having a black end and a white end. Constellations are generalizations of bipartite maps with more colors (see [Section 2](#) for precise definitions).

Map enumeration has been an important research topic for many years now, going back to Tutte [\[22\]](#) with planar maps. He used analytic techniques on generating functions, and later on, Schaeffer enumerated planar maps bijectively [\[21\]](#), with many generalizations (see for instance [\[5, 4, 1\]](#)). The enumeration of maps was extended to other models : for instance, asymptotic formulas were obtained by Bender and Canfield [\[3\]](#) for maps of higher genus, by Gao [\[12\]](#) for maps with prescribed degrees, and Chapuy [\[7\]](#) for constellations. Another way to count maps is to see them as factorizations of permutations and to use algebraic properties of \mathfrak{S}_n . In particular, maps fit in the more general context of weighted Hurwitz numbers (see e.g. [\[2\]](#)). Their generating functions satisfy integrable hierarchies of PDEs that arose from mathematical physics, namely the KP and 2-Toda hierarchies (a good introduction can be found in [\[17\]](#)).

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The first numbers that were studied from the point of view of integrable hierarchies were Hurwitz numbers, that enumerate ramified coverings of the sphere. Pandharipande conjectured a recurrence formula for those numbers [20], which was proven by Okounkov [18] and later simplified by Dubrovin, Yang and Zagier [9]. Later, recurrence formulas for maps were found, starting with Goulden and Jackson for triangulations [14]. They were followed by Carrell and Chapuy for general maps [6], and Kazarian and Zograf for bipartite maps [15]. All these works start from the fact that an underlying generating function is a "tau function" of an integrable hierarchy, and then use ad-hoc techniques to obtain explicit recurrence formulas. The generality of this second step is not well understood. In particular, although the methods mentioned in the previous page enable to control face degrees, this is not the case of the efficient recurrences of [9, 14, 6, 15]. This raises the question of obtaining similar formulas for other models of maps.

Contributions of this article: We show that similar techniques as in [18, 9] apply more generally in the context of maps, and we derive recurrence formulas for bipartite maps with prescribed degrees, allowing us in particular to derive a formula for bipartite $2k$ -angulations. We also find recurrence formulas for constellations. These formulas are, up to our knowledge, the simplest and fastest way to calculate those numbers (in all models, it takes $O(n^2g^3)$ arithmetic operations to calculate the coefficient for n edges and genus g , see [Remark 1](#)). We insist on the fact that the underlying "tau function" structure is well-known, the contribution of this paper is to show that it can be used to control face degrees and count constellations.

Structure of the paper: In [Section 2](#), we will give precise definitions and state our main results. The rest of the paper presents the main steps of the proof. The first part of the proof is common to all models : we introduce the "tau function" τ , a certain generating function for constellations. This function, along with some auxiliary functions τ_n , classically satisfies a set of differential equations called the 2-Toda hierarchy. Our first contribution, inspired by [18], is to link τ to the τ_n and derive an equation involving τ only ([Proposition 6](#)). This will be presented in [Section 3](#). From this equation, specialized to the model we wish for (bipartite maps or constellations), we perform a few combinatorial operations (that are specific to the model, similarly as in [6, 14, 15]) to obtain our formulas. We will present this in details for bipartite maps in [Section 4](#), and we briefly mention the case of constellations.

2 Definitions and main results

Definition 1. A map M is the data of a connected multigraph (multiple edges and loops are allowed) G (called the underlying graph) embedded in a compact oriented surface S , such that $S \setminus G$ is homeomorphic to a collection of disks. The connected components of

$S \setminus G$ are called the *faces*. The *genus* g of M is the genus of S (the number of "handles" in S). M is defined up to orientation-preserving homeomorphism. A *bipartite map* is a map with two types of vertices (black or white), such that each edge connects two vertices of different colors. A bipartite map is said to be *rooted* if a particular edge is distinguished.

An *m-constellation* is a particular kind of map with special vertices : colored vertices, carrying a "color" between 1 and m , and star vertices. Each edge connects a star vertex to a colored vertex. A star vertex has degree m , and its neighbors have color $1, 2, \dots, m$ in the clockwise cyclic order. Bipartite maps are in bijection with 2-constellations, since each star vertex and its two adjacent edges can be merged into a single edge connecting a black and a white vertex. A constellation is said to be rooted if a particular star vertex is distinguished. A constellation with n star vertices is said to be *labeled* if each star vertex carries a different label between 1 and n . Since rooting kills all possible automorphisms, there is a $(n - 1)!$ -to-1 correspondence between labeled and rooted constellations with n star vertices. From now on, we will only consider rooted objects unless stated otherwise.

Some basic, well-known, properties of maps and constellations will be useful later.

Proposition 1. *Labeled (non-necessarily connected) m-constellations with n star vertices are in bijection with (m + 1)-uples $(\sigma_1, \sigma_2, \dots, \sigma_m, \phi)$ of permutations of \mathfrak{S}_n such that $\sigma_1 \cdot \sigma_2 \cdot \dots \cdot \sigma_m = \phi$. The permutation σ_i represents the vertices of color i : each vertex is a cycle of σ_i , and the elements of the cycle represent the neighboring star vertices, in that cyclic order. The permutation ϕ encodes the faces. See Figure 1 for an example.*

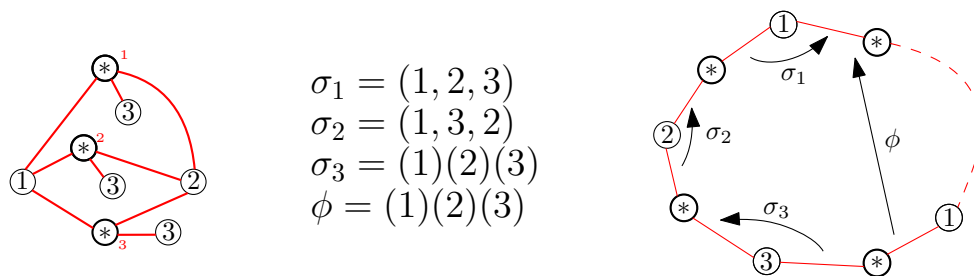


Figure 1: Left : a (labeled) 3-constellation (of genus 0) and the corresponding permutations, right : the permutation ϕ , whose cycles describe the faces.

Our main results are the following theorems :

Theorem 2. *The number $A_{g,n}^{(k)}$ of bipartite $2k$ -angulations of genus g with n faces satisfies the following recurrence formula :*

$$\binom{kn+1}{2} A_{g,n}^{(k)} = \sum_{\substack{n_1+n_2=n \\ n_1, n_2 \geq 1 \\ g_1+g_2+g^*=g}} (kn_1+1) \binom{(k-1)n_2+2-2g_2}{2g^*+2} A_{g_1, n_1}^{(k)} A_{g_2, n_2}^{(k)} \\ + \sum_{g^* \geq 0} \binom{(k-1)n+2-2(g-g^*)}{2g^*+2} A_{g-g^*, n}^{(k)}.$$

Theorem 2 is actually a particular case of the following theorem for bipartite maps with prescribed degrees :

Theorem 3. *The number $B_g(\mathbf{f})$ of bipartite maps of genus g with f_i faces of degree $2i$ (for $\mathbf{f} = (f_1, f_2, \dots)$) satisfies :*

$$\binom{n+1}{2} B_g(\mathbf{f}) = \sum_{\substack{s+t=f \\ s, t \neq 0 \\ g_1+g_2+g^*=g}} (1+n_1) \binom{v_2}{2g^*+2} B_{g_1}(\mathbf{s}) B_{g_2}(\mathbf{t}) + \sum_{g^* \geq 0} \binom{v+2g^*}{2g^*+2} B_{g-g^*}(\mathbf{f}) \quad (2.1)$$

where $n = \sum_i i f_i$, $n_1 = \sum_i i s_i$, $v = 2 - 2g + n - \sum_i f_i$, $v_2 = 2 - 2g_2 + n_2 - \sum_i t_i$ and $n_2 = \sum_i i t_i$ (the n 's count edges, the v 's count vertices, in accordance with the Euler formula).

Theorem 4. *The numbers $C_{g,n}^{(m)}$ of m -constellations of genus g with n star vertices satisfy the following recurrence formula :*

$$\binom{n}{2} C_{g,n}^{(m)} = \sum_{\substack{n_1+n_2=n \\ n_1, n_2 \geq 1 \\ g=g_1+g_2+g^*}} n_1 \binom{(m-1)n_2+2-2g_2}{2g^*+2} C_{g_1, n_1}^{(m)} C_{g_2, n_2}^{(m)}.$$

Remark 1. **Theorem 3** allows to compute the number of maps with prescribed degrees way faster than the usual Tutte-Lehman-Walsh approach [23, 3, 12] or the topological recursion (see e.g. [10]), especially for higher genus (because these methods require counting maps with up to g boundaries to enumerate maps of genus g). It can also be specialized to maps with bounded face degrees (contrarily to the Tutte equation).

We observe that **Theorem 4** applies to bipartite maps (for $m = 2$). It is also possible to obtain an analogue of **Theorem 3** (with prescribed degrees) for m -constellations with $m \geq 3$, but the formula is more complicated and is not quadratic anymore (we do not include it in this short version). We give a brief explanation of that fact in **Section 4.2**.

We can also derive a similar-looking formula for monotone simple Hurwitz numbers, see (5.1).

3 Generating functions for constellations

Our goal is to prove [Proposition 6](#). In order to do this, we first have to introduce Maya diagrams and the semi infinite wedge space.

3.1 The semi infinite wedge space

We give some definitions, mostly following the notations of the appendix in [19] :

Definition 2. A *Maya diagram* is a decoration of $\mathbb{Z} + \frac{1}{2}$ with a particle or an antiparticle at each position, such that for some n_1, n_2 there are only particles at positions $< n_1$ and only antiparticles at positions $> n_2$. The semi infinite wedge space Λ^{∞} is the vector space whose orthonormal basis elements are the Maya diagrams.

For any $k \in \mathbb{Z} + \frac{1}{2}$, we define the *fermion operators* ψ_k and ψ_k^* . For each Maya diagram \mathbf{m} , we define :

$$\psi_k \mathbf{m} = \begin{cases} 0 & \text{if } \mathbf{m} \text{ has a particle in position } k \\ \tilde{\mathbf{m}} & \text{otherwise} \end{cases}$$

$$\psi_k^* \mathbf{m} = \begin{cases} 0 & \text{if } \mathbf{m} \text{ has an antiparticle in position } k \\ \bar{\mathbf{m}} & \text{otherwise} \end{cases}$$

where $\tilde{\mathbf{m}}$ is the same as \mathbf{m} except there is a particle in position k , and $\bar{\mathbf{m}}$ is the same as \mathbf{m} except there is an antiparticle in position k .

We can now define the *boson operators* : for all $n \in \mathbb{Z} \setminus \{0\}$, let

$$\alpha_n = \sum_{k \in \mathbb{Z} + \frac{1}{2}} \psi_{k-n} \psi_k^*.$$

Finally, the two *vertex operators* are

$$\Gamma_{\pm}(\mathbf{p}) = \exp\left(\sum_{n=1}^{\infty} \frac{p_n}{n} \alpha_{\pm n}\right).$$

Definition 3. We define the *normally ordered products*

$$: \psi_k \psi_k^* : := \begin{cases} \psi_k \psi_k^* & \text{if } k > 0 \\ -\psi_k^* \psi_k & \text{if } k < 0 \end{cases}$$

and the *charge* and *energy operators* :

$$C = \sum_{k \in \mathbb{Z} + \frac{1}{2}} : \psi_k \psi_k^* : \quad H = \sum_{k \in \mathbb{Z} + \frac{1}{2}} k : \psi_k \psi_k^* :$$

The Maya diagrams are eigenvectors of C . The eigenvalue of a Maya diagram \mathbf{m} is the number of particles in positive position minus the number of antiparticles in negative position. We call this number the *charge* of \mathbf{m} . We introduce the translation operator R : for any \mathbf{m} , $R\mathbf{m}$ has a particle in position $k + 1$ if and only if \mathbf{m} has a particle in position k . Note that if the charge of \mathbf{m} is c , the charge of $R\mathbf{m}$ is $c + 1$.

There is a bijection between Maya diagrams of charge 0 and partitions, as depicted in [Figure 2](#). Thus, any Maya diagram \mathbf{m} can be encoded by its charge c and a partition λ (that corresponds to the Maya diagram $R^{-c}\mathbf{m}$).

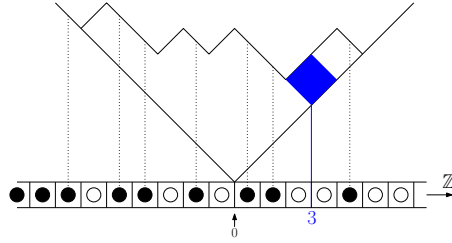


Figure 2: A maya diagram of charge 0 and its corresponding partition (above it, presented as a rotated Young diagram). In blue, a box and its content (the abscissa of the projection of the center of the box on \mathbb{Z})

We will use the bracket notation, and denote the Maya diagram corresponding to the empty partition by $|\emptyset\rangle$, and set $|\emptyset_n\rangle = R^n|\emptyset\rangle$. We will also set $|\lambda\rangle$ to be the Maya diagram of charge 0 corresponding to the partition λ . In particular, $H|\lambda\rangle = |\lambda||\lambda\rangle$, where $|\lambda|$ is the number of boxes in λ .

3.2 Constellations and content products

Definition 4. Fix integers r, n and g , fix λ, μ two partitions of n . Let $W_n^{\lambda, \mu}(l_1, l_2, \dots, l_r)$ be the number of $(r + 2)$ -uples of permutations $(\sigma_1, \sigma_2, \dots, \sigma_r, \sigma_\lambda, \sigma_\mu)$ of \mathfrak{S}_n such that $\sigma_1 \cdot \sigma_2 \cdot \dots \cdot \sigma_r = \sigma_\lambda \sigma_\mu$ and σ_i has l_i cycles, and $\sigma_\lambda, \sigma_\mu$ have respective cycle types λ and μ . The $W_n^{\lambda, \mu}$ enumerate (labeled, non-necessarily connected) constellations, in accordance with [Proposition 1](#). Let τ be the associated generating function (that implicitly depends on r):

$$\tau(z, \mathbf{p}, \mathbf{q}, (u_j)) = \sum_{\substack{n \geq 0 \\ |\mu| = |\lambda| = n \\ l_i \geq 1 \forall i}} \frac{z^n}{n!} \prod_{i=1}^r u_i^{n-l_i} p_\lambda q_\mu W_n^{\lambda, \mu}(l_1, l_2, \dots, l_r)$$

It is a classical result (under different forms and variants, see for instance [\[14, 18\]](#)) that the function τ can be expressed in terms of elements and operators of $\Lambda^{\frac{\infty}{2}}$:

Lemma 5 (Classical).

$$\tau(z, \mathbf{p}, \mathbf{q}, (u_j)) = \langle \emptyset | \Gamma_+(\mathbf{p}) z^H \Lambda \Gamma_-(\mathbf{q}) | \emptyset \rangle \quad (3.1)$$

with

$$F(u) = \sum_{k>0} \sum_{i=0}^{k-1/2} \log(1+ui) \psi_k \psi_k^* + \sum_{k<0} \sum_{i=0}^{-k-1/2} \log(1-ui) \psi_k^* \psi_k$$

and $\Lambda = \prod_{j=1}^r \exp(F(u_j))$.

Proof. For reasons of space, we only include the main ideas of the proof.

First, we have

$$F(u) |v\rangle = \left(\sum_{\square \in v} \log(1+uc(\square)) \right) |v\rangle$$

where the $c(\square)$ are the contents of the partition v (see [Figure 2](#)). It can be shown using the Jacobi-Trudi rule (see e.g. [\[19\]](#)) that

$$\Gamma_-(\mathbf{q}) | \emptyset \rangle = \sum_v s_v(\mathbf{q}) |v\rangle \quad \text{and} \quad \langle \emptyset | \Gamma_+(\mathbf{p}) = \sum_v s_v(\mathbf{p}) \langle v|.$$

Thus the RHS of [\(3.1\)](#) can be rewritten as :

$$\sum_{\substack{n>0 \\ |v|=n}} z^n \prod_{j=1}^r \prod_{\square \in v} (1+u_j c(\square)) s_v(\mathbf{p}) s_v(\mathbf{q}).$$

Here we recognize the "content product form" of the generating function of constellations, as in [\[14\]](#), which concludes the proof. \square

We introduce the auxiliary functions τ_n , for $n \in \mathbb{Z}$:

$$\tau_n = \langle \emptyset_n | \Gamma_+(\mathbf{p}) z^H \Lambda \Gamma_-(\mathbf{q}) | \emptyset_n \rangle.$$

We have $\tau = \tau_0$. The previous lemma, along with classical considerations (see for instance Section 2.6 in [\[18\]](#)), imply that the τ_n satisfy an infinite set of equations, the 2-Toda hierarchy. In particular, the following equation holds :

$$\frac{\partial^2}{\partial p_1 \partial q_1} \log \tau_0 = \frac{\tau_1 \tau_{-1}}{\tau_0^2}. \quad (3.2)$$

So far, the content presented was classical. Our first main contribution is to transform the previous equation into an equation implying τ only.

Proposition 6. *The function τ introduced in [Definition 4](#) satisfies*

$$\frac{\partial^2}{\partial p_1 \partial q_1} \log \tau = z \frac{\tau(z \cdot \prod_{j=1}^r (1+u_j), \mathbf{p}, \mathbf{q}, (\frac{u_j}{1+u_j})) \tau(z \cdot \prod_{j=1}^r (1-u_j), \mathbf{p}, \mathbf{q}, (\frac{u_j}{1-u_j}))}{\tau^2}. \quad (3.3)$$

Proof. It is easily verified that R commutes with $\Gamma_+(\mathbf{p})$ and $\Gamma_-(\mathbf{q})$, as well as $R^{-n}HR^n = H + nC + \frac{n^2}{2}$. Thus,

$$\tau_n = \langle \emptyset | R^{-n} \Gamma_+(\mathbf{p}) z^H \Lambda \Gamma_-(\mathbf{q}) R^n | \emptyset \rangle = \langle \emptyset | \Gamma_+(\mathbf{p}) z^{H+nC+\frac{n^2}{2}} R^{-n} \Lambda R^n \Gamma_-(\mathbf{q}) | \emptyset \rangle.$$

But, by a careful change of indices,

$$\begin{aligned} RFR^{-1} &= \sum_{k>0} \sum_{i=0}^{k-1/2} \log(1+ui) \psi_{k+1} \psi_{k+1}^* + \sum_{k<0} \sum_{i=0}^{-k-1/2} \log(1-ui) \psi_{k+1}^* \psi_{k+1} \\ &= \sum_{k>0} \sum_{i=1}^{k-1/2} \log(1-u+ui) \psi_k \psi_k^* + \sum_{k<0} \sum_{i=1}^{-k-1/2} \log(1-u-ui) \psi_k^* \psi_k \\ &= F\left(\frac{u}{1-u}\right) + (H-C/2) \log(1-u). \end{aligned}$$

Similarly,

$$R^{-1}FR = F\left(\frac{u}{1+u}\right) + (H+C/2) \log(1+u).$$

Thus, we get :

$$\tau_{\pm 1}(z, \mathbf{p}, \mathbf{q}, (u_j)) = z^{1/2} \tau\left(z \cdot \prod_{j=1}^r (1 \pm u_j), \mathbf{p}, \mathbf{q}, \left(\frac{u_j}{1 \pm u_j}\right)\right),$$

and the result follows by plugging this into (3.2). \square

Remark 2. The idea of expressing $\tau_{\pm 1}$ in terms of τ by calculating $R^{\mp 1} \Lambda R^{\pm 1}$ is inspired by the calculation performed in [18], Section 2.7.

4 Recurrence formulas

4.1 Bipartite maps

In this section we study bipartite maps, therefore we work with $r = 2$. The objects we are interested in are connected, thus, in accordance with the "exp-log principle" (see for instance [11]), we need to consider $H = \log(\tau)$. Let $H^{\lambda, \mu}$ be its coefficients in the series expansion, namely

$$H = \sum_{n \geq 0, |\mu| = |\lambda| = n} \frac{z^n}{n!} H^{\lambda, \mu} p_{\lambda} q_{\mu}.$$

Note that $H^{\lambda, \mu}$ is a polynomial in u_1 and u_2 . We will set $u_1 = u_2$ (which is equivalent to counting vertices regardless of their color) and "get rid" of the \mathbf{q} variables. More

precisely, let $H_n^\lambda = H^{\lambda, 1^n} |_{u_1=u_2=u}$ and

$$H^*(z, \mathbf{p}, u) = H |_{q_i=\delta_{i,1}, u_1=u_2=u} = \sum_{n \geq 0, |\lambda|=n} H_n^\lambda \frac{z^n}{n!} p_\lambda.$$

By **Proposition 1**, H_n^λ counts (connected, labeled) bipartite maps with n edges, with λ describing the face type and a weight u per vertex.

The next step, inspired by the case of classical Hurwitz numbers [9], is to obtain a quadratic (differential) equation for H^* from (3.3).

Lemma 7.

$$D^2 H_1^* - D H_1^* = D H_1^* \left(D H^*(z(1+u)^2, \mathbf{p}, \frac{u}{1+u}) + D H^*(z(1-u)^2, \mathbf{p}, \frac{u}{1-u}) - 2 D H^* \right) \quad (4.1)$$

with $H_1^* = \frac{\partial}{\partial p_1} H^*$ and $D = z \frac{\partial}{\partial z}$.

Proof. We have

$$\frac{\partial^2}{\partial p_1 \partial q_1} \log \tau \Big|_{q_i=\delta_{i,1}, u_1=u_2=u} = \left(\sum_{n>0} H_n^\lambda \frac{z^n}{n!} n p_\lambda q_1^{n-1} \right) \Big|_{q_1=1} = z \frac{\partial}{\partial z} H_1^*.$$

Hence (3.3) now reads

$$D H_1^*(z, \mathbf{p}, u) = z \exp \left(H^*(z(1+u)^2, \mathbf{p}, \frac{u}{1+u}) + H^*(z(1-u)^2, \mathbf{p}, \frac{u}{1-u}) - 2 H^* \right).$$

Applying D to the previous equation, and eliminating the exponential term between the two equations, one gets the result. \square

We will now interpret H^* and H_1^* in (4.1) in terms of bipartite maps, using a local combinatorial operation that is specific to bipartite maps. This step is conceptually similar to what was done in [14, 6, 15].

Lemma 8. Let $B(z, \mathbf{p}, u)$ be the ordinary generating function of connected rooted bipartite maps, and $B_g(\mathbf{f})$ be the number of bipartite maps of genus g with f_i faces of degree $2i$, such that

$$B = \sum_{g, \mathbf{f}} z^d u^{2n-v} \prod_{i \geq 1} p_i^{f_i} B_g(\mathbf{f}).$$

with $n = \sum_i i f_i$ and $v - n + \sum_i f_i = 2 - 2g$ (Euler formula). B satisfies the following equation :

$$(D+1)DB = (u^{-2} + (D+1)B) \left(B(z(1+u)^2, \mathbf{p}, \frac{u}{1+u}) + B(z(1-u)^2, \mathbf{p}, \frac{u}{1-u}) - 2B \right). \quad (4.2)$$

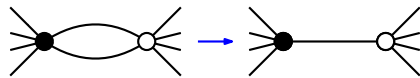


Figure 3: Contraction of a digon.

Proof. H^* is the exponential generating function of labeled connected bipartite maps, where u records the number of vertices, z records edges, and p_i records faces of degree $2i$. We directly have $B = DH^*$, as there is an $(n-1)!$ -to-1 correspondence between labeled and rooted bipartite maps.

The function $p_1 H_1^*$ would be the EGF of the same objects, with a marked digon (face of degree 2). There is a $n!$ -to-1 correspondence between labeled and unlabeled-unrooted bipartite maps with a marked digon and n edges. A marked digon can be contracted into a root edge (see [Figure 3](#)) except when the bipartite map is just one edge, thus $H_1^* = z + u^2 z B$ (the $u^2 z$ factor comes from the fact that we lose an edge when we contract the digon, and the z term is the case where we cannot contract the digon). □

We are finally ready to prove [Theorem 3](#).

Proof of [Theorem 3](#). We look at the factor $B(z(1+u)^2, \mathbf{p}, \frac{u}{1+u}) + B(z(1-u)^2, \mathbf{p}, \frac{u}{1-u}) - 2B$ in [\(4.2\)](#). The coefficient of $z^n \prod_{i \geq 1} p_i^{f_i}$ in it is :

$$\sum_{v>0} B_g(\mathbf{f}) u^{2n-v} ((1+u)^v + (1-u)^v - 2) = \sum_{v>0} B_g(\mathbf{f}) u^{2n-v} \left(2 \sum_{0 < k \leq \frac{v}{2}} u^{2k} \binom{v}{2k} \right).$$

Extracting coefficients in [\(4.2\)](#), one gets the result. □

Remark 3. If we don't specialize $u_1 = u_2 = u$ as we did at the beginning of this subsection, we can obtain formulas for more precise coefficients of bipartite maps (namely knowing the number of vertices of each color). This will appear in the long version.

4.2 Constellations

The proof of [Theorem 4](#) is similar to the proof of [Theorem 3](#) that was presented in the previous section : for m -constellations, we consider our function τ with $r = m + 1$. However, we now "get rid" of both the \mathbf{p} variables and the \mathbf{q} variables by specializing $p_i = q_i = \delta_{i,1}$, hence we do not control over the degrees of the faces.

As mentioned in the introduction, one can also study the effect of the specialization $q_i = \delta_{i,1}$ (leaving the p_i non specialized) in the case of constellations with $m > 2$. However the operation of contracting an m -gon (that can be defined similarly as the contraction of a digon presented in [Figure 3](#)) might disconnect the map, and the formula we obtain is more complicated (it will be included in a full version).

5 Conclusion

The coefficients in our recurrence formulas have a combinatorial flavor. It is a natural question to ask for a bijective proof of these formulas. However, the bijective interpretation of formulas derived from the KP/2-Toda hierarchies is still a widely open question, as bijections have only been found for certain formulas, in the particular cases of one-faced [8] and planar maps [16].

Using the same technique, we can also obtain recurrence formulas for similar numbers called monotone simple (i.e. unramified at 0 and ∞) Hurwitz numbers, introduced in [13], that are also particular cases of weighted Hurwitz numbers. These numbers $\vec{H}_{g,n}$ (of genus g , in \mathfrak{S}_n) satisfy the following recurrence formula :

$$(n^3 - n^2)\vec{H}_{g,n} = 2 \sum_{\substack{n_1+n_2=n \\ g^* \geq 0 \\ g_1+g_2+g^*=g}} n_1^2 n_2 \binom{2n_2 + 2g_2 + 2g^* - 1}{2g^* + 2} \vec{H}_{g_1, n_1} \vec{H}_{g_2, n_2}. \quad (5.1)$$

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